

1. 40 points.

(a) The van der Waals b coefficient for CCl_2F_2 is $0.0998 \text{ L mol}^{-1}$. Estimate R_{LJ} . **Solution:**

$$b \approx \frac{2\mathcal{N}_A\pi R_{\text{LJ}}^3}{3} \quad R_{\text{LJ}} \approx \left(\frac{3b}{2\mathcal{N}_A\pi}\right)^{1/3} = \boxed{4.3 \text{ \AA}}$$

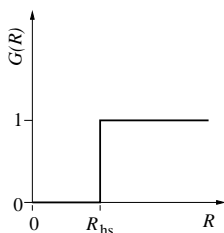
(b) Give the value of each probability for a rotating linear molecule:

i. $\lim_{T \rightarrow 0} \mathcal{P}(J=0) = \boxed{1}$ (all molecules in ground state.)

ii. $\lim_{T \rightarrow 0} \mathcal{P}(J=4) = \boxed{0}$ (all molecules in ground state.)

iii. $\lim_{T \rightarrow \infty} \mathcal{P}(J=4) = \boxed{0}$ (molecules equally spread over all states.)

(c) Sketch the pair correlation function for a hard sphere potential.



(d) Identify each of the following particles as a fermion (“F”) or boson (“B”).

i. electron F

ii. ^{235}U nucleus F

iii. ^{235}U neutral atom B

iv. $^{19}\text{F}^-$ ion F

v. $^1\text{H}_2$ molecule B

2. Calculate the pressure in bar of 1210 mol krypton in a 150.0 L container at 298 K, if $a = 2.325 \text{ L}^2 \text{ mol}^{-2}$ and $b = 0.0396 \text{ L mol}^{-1}$. **Solution:**

$$\begin{aligned} RT &= \left(P + \frac{a}{V_m^2}\right)(V_m - b) \\ P &= \frac{RT}{V_m - b} - \frac{a}{V_m^2} \\ &= \frac{(0.083145 \text{ bar L K}^{-1} \text{ mol}^{-1})(298 \text{ K})}{(150.0 \text{ L}/1210 \text{ mol}) - 0.0396 \text{ L mol}^{-1}} - \frac{2.325 \text{ L}^2 \text{ mol}^{-2}}{(150.0 \text{ L}/1210 \text{ mol})^2} \\ &= \boxed{142 \text{ bar}} \end{aligned}$$

The ideal gas pressure would be 200 bar.

3. Find the integral approximation to the vibrational partition function for a degenerate bend in a linear molecule, where $E_{\text{vib}} = \omega_e v$ and $g_{\text{vib}} = v + 1$. **Solution:** For the integral approximation, we replace the sum by an integral and multiply by the volume element dv , which effectively has

a value of 1:

$$\begin{aligned}
 q_{\text{vib}} &= \sum_0^{\infty} (v+1) e^{-\omega_e v / (k_B T)} \approx \int_0^{\infty} (v+1) e^{-\omega_e v / (k_B T)} dv \\
 &= \int_0^{\infty} v e^{-\omega_e v / (k_B T)} dv + \int_0^{\infty} e^{-\omega_e v / (k_B T)} dv \\
 &= \frac{1}{\omega_e^2 / (k_B T)^2} + \frac{1}{\omega_e / (k_B T)} = \frac{(k_B T)^2 + \omega_e (k_B T)}{\omega_e^2}.
 \end{aligned}$$

In the limit that $k_B T \gg \omega_e$, this simplifies to $(k_B T / \omega_e)^2$, which is the same result predicted in the high-temperature limit from multiplying together the two partition functions for each bending mode, $q_{\text{vib}}(1)q_{\text{vib}}(2)$.

4. For a crystal where each vibrational mode can absorb a finite amount of energy per mode,

$$\epsilon(T) = \frac{\epsilon_0 k_B T}{\epsilon_0 + k_B T},$$

(a) find an equation for C_{Vm} ; (b) find $\lim_{T \rightarrow 0} C_{Vm}$; (c) find $\lim_{T \rightarrow \infty} C_{Vm}$. **Solution:** There are still $3N$ vibrational modes, as in the Einstein theory, so the molar energy will be $E_m = 3N_A \epsilon(T)$. To get the heat capacity, we'll take the derivative with respect to temperature:

$$\begin{aligned}
 C_{Vm} &= \left(\frac{\partial E_m}{\partial T} \right)_V = \frac{\partial}{\partial T} \left(\frac{3N_A \epsilon_0 k_B T}{\epsilon_0 + k_B T} \right) \\
 &= 3N_A k_B \epsilon_0 \left[\frac{1}{\epsilon_0 + k_B T} - \frac{T}{(\epsilon_0 + k_B T)^2} \right] = \frac{3R \epsilon_0}{\epsilon_0 + k_B T} \left[1 - \frac{T}{\epsilon_0 + k_B T} \right] \\
 \lim_{T \rightarrow 0} C_{Vm} &= \frac{3R \epsilon_0}{\epsilon_0 + 0} \left[1 - \frac{0}{\epsilon_0 + 0} \right] = \boxed{3R} \\
 \lim_{T \rightarrow \infty} C_{Vm} &= \lim_{T \rightarrow \infty} \frac{3R \epsilon_0}{k_B T} \left[1 - \frac{T}{k_B T} \right] = \boxed{0}.
 \end{aligned}$$